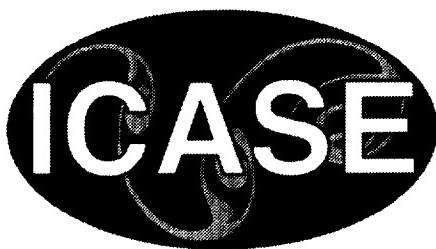


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NUMERICAL STUDY OF INTERACTION OF A VORTICAL DENSITY INHOMOGENEITY WITH SHOCK AND EXPANSION WAVES

A. POVITSKY * AND D. OFENGEIM †

Abstract. We studied the interaction of a vortical density inhomogeneity (VDI) with shock and expansion waves. We call the VDI the region of concentrated vorticity (vortex) with a density different from that of ambiance. Non-parallel directions of the density gradient normal to the VDI surface and the pressure gradient across a shock wave results in an additional vorticity. The roll-up of the initial round VDI towards a non-symmetrical shape is studied numerically.

Numerical modeling of this interaction is performed by a 2-D Euler code. The use of an adaptive unstructured numerical grid makes it possible to obtain high accuracy and capture regions of induced vorticity with a moderate overall number of mesh points. For the validation of the code, the computational results are compared with available experimental results and good agreement is obtained.

The interaction of the VDI with a propagating shock wave is studied for a range of initial and induced circulations and obtained flow patterns are presented. The splitting of the VDI develops into the formation of a non-symmetrical vortex pair and not in a set of vortices.

A method for the analytical computation of an overall induced circulation Γ_1 as a result of the interaction of a moving VDI with a number of waves is proposed. Simplified, approximated, expressions for Γ_1 are derived and their accuracy is discussed.

The splitting of the VDI passing through the Prandtl-Meyer expansion wave is studied numerically. The obtained VDI patterns are compared to those for the interaction of the VDI with a propagating shock wave for the same values of initial and induced circulations. These patterns have similar shapes for corresponding time moments.

Key words. vortices, unstructured grids, automatic mesh refinement, Euler equations, vortical density inhomogeneity, shock wave, expansion wave

Subject classification. Fluid Mechanics

1. Introduction. A common situation in compressible flows is that regions of concentrated vorticity (vortex) are also the regions of inhomogeneous density, i.e., different density than ambient. This vortical density inhomogeneity (VDI) interacts with either a shock or an expansion wave, i.e., with regions of pressure gradients. Current interest in this type of interaction is motivated to promote rapid mixing and combustion of hydrogen and hydrocarbon fuels for supersonic combustor applications and problems of direct numerical simulation of compressible turbulence. This research is a numerical investigation of this type of interaction.

Previous investigations of this type of interaction were concerned with the interaction of either a density inhomogeneity (initial vorticity is zero) [1] or a "pure" vortex [2] (without initial density inhomogeneity) with a shock wave.

Non-parallel directions of the density gradient normal to a VDI surface and the pressure gradient across

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a shock wave cause an additional vorticity (denoted also as induced vorticity). This vorticity causes the VDI to roll up into a non-symmetrical shape and to evolve towards a set of vortices of a finite core size. Detailed experiments were carried out by Jacobs [3] to quantify the mixing induced by the interaction of a shock wave with a cylindrical volume of a light gas (helium). Yang et al. [4] obtained an approximate analytical expression for this type of interaction.

A "pure" vortex interaction with a shock wave has been studied in various articles including the detailed numerical research reported in [5]. The experimental research [6] has shown that the vortex is somewhat compressed by the shock, while it generally retains the same configuration.

However, the interaction of the VDI with a shock wave as well as splitting of the VDI due to its passage through regions of pressure gradients, like expansion waves, has not been studied previously.

Generally, the presence of an induced and initial circulation can lead to the roll-up of the VDI into non-regular shapes, commonly including a number of vortices of different intensity. Our computations show that vortices of the same sign merge into a single vortex and VDI shapes transform into a non-symmetrical vortex pair of finite core size.

Our computations of the VDI interaction with an expansion wave show that in spite of obvious differences between interactions of VDI with shock and expansion waves, similar VDI shapes are obtained for the same values of induced and initial circulation. A correct estimation of induced circulation for the interaction with an expansion wave should be done. In this article, we derive it for the Prandtl-Meyer expansion wave.

Special numerical techniques were used by various authors for capturing the roll-up of a gas bubble moving with the flowfield after the shock-bubble interaction. Yang et al [4] use a regular numerical grid and rectangular computational domain moving in time to track the developing vortex pair. Quirk and Karni [7] use the Adaptive Mesh Refinement (AMR) algorithm. The AMR algorithm employs hierarchical system of rectangular grids. A local solution at a previous time step is used to determine automatically where refinement is needed to resolve small-scale phenomena.

An unstructured grid Euler solver [8]-[11] with a local refinement based on a density gradient is used in this research. Its advantage is an automatic capture of thin regions of induced vorticity at the VDI boundary.

In section 2, governing equations and the computational method used are described. In section 3, we investigate the phenomena of the VDI splitting due to the interaction with a propagating shock wave and obtain dynamics of VDI patterns for a range of induced and initial circulation. In the last section, we derive analytical estimations for induced circulation due to the VDI interaction with the Prandtl-Meyer expansion wave and obtain VDI shapes by numerical computations.

2. Governing equations and numerical method. For a short limited time (a few milliseconds) viscous effects are minor and a non-viscous assumption is used to describe the phenomenon [4], [7].

The two-dimensional unsteady Euler equations of gas dynamics describing an inviscid compressible flow in conservative form can be written as:

$$(1) \quad \mathbf{u}_t + \mathbf{f}_x + \mathbf{g}_y = 0.$$

The vectors in equation (1) are

$$(2) \quad \mathbf{u} = \begin{pmatrix} \rho \\ \rho u \\ \rho v \\ e \end{pmatrix}, \quad \mathbf{f} = \begin{pmatrix} \rho u \\ \rho u^2 + p \\ \rho uv \\ (e + p)u \end{pmatrix}, \quad \mathbf{g} = \begin{pmatrix} \rho v \\ \rho uv \\ \rho u^2 + p \\ (e + p)v \end{pmatrix},$$

where ρ is density, p -static pressure, (u, v) – velocity in Cartesian coordinates (x, y) , and e is the total energy, related to the other variables by an equation of state which for a perfect gas is

$$(3) \quad e = \frac{p}{(\gamma - 1)} + \frac{1}{2} \rho(u^2 + v^2)$$

where $\gamma = 1.4$ ratio of specific heats.

The inviscid calculations have been carried out using an Euler code [10]. The code combines advantages of the total variation diminishing (TVD) finite-volume schemes, unstructured grids and an unsteady local adaptation of the grid to the features of the flow. The second-order Godunov-type scheme modified for transient flows and unstructured grids by Fursenko et al. [9], was employed to solve the Euler equations. This scheme is based on the TVD approach combined with second-order upwind differences and a solution of the Riemann problem to obtain inter-zone fluxes. More detailed description of the numerical method used in the Euler code can be found in [8]-[9].

The unstructured grids are composed of triangular area elements, that provide a more isotropic spatial discretization as compared to traditional rectangular grids [8]. In order to generate an unstructured grid in the computational domain the 2-D unstructured grid generator created by Voinovich and Galyukov [11] was used. Fragments of the unstructured grid for a typical VDI interaction with a propagating shock wave are shown in Figure 1a-b. The distance between the neighboring grid points is 2^k times smaller than the initial distance. Computations are done with the maximum k equal to 5.

The graphical presentation of numerical results has been done by the grid data VIGIE code [12] that allows to visualize an unstructured grid computed fluid flow.

There are no available experiments for interactions with the VDI; therefore, results of an experimental research for the interaction of a moving shock wave with density inhomogeneity [3] are used for the code verification. Experiments were carried out for the interaction of a shock wave ($M=1.09$) with a cylindrical volume of gas (helium) that is lighter than ambiance (air). Computational results are presented in Fig 2. There is a good agreement with experimental results [3] at corresponding time moments.

3. Interaction between vortical density inhomogeneity and a propagating shock wave. A computational domain for the interaction between the VDI and a straight propagating shock wave is shown in Figure 3.

The initial conditions behind a shock wave are calculated according to the Rankine-Hugoniot relations; ahead of the shock wave ambient conditions are given. The shock wave proceeds into ambient conditions and interacts with a vortex superimposed on the ambient state.

The initial density ratio ($D = \rho_l / \rho_h$) between a light VDI material (ρ_l) and a heavy ambient air (ρ_h) and the velocity ratio ($W = U_c / U_b$) between a maximum tangential velocity (U_c) and a velocity behind the shock wave (U_b) are governing parameters that determine the dynamics of the VDI shape after an interaction with a propagating shock wave.

The 2-D vortex model [5] used for the VDI consists of two regions of the vortical flow: an inner core region (R_c) and a surrounding region (R_o) where the velocity gradually reduces to zero (see Fig. 3a). In the numerical experiments carried out the linear tangential velocity profile in the VDI core region is considered and is represented by

$$(4) \quad U_\theta(r) = U_c \frac{r}{R_c}; r < R_c,$$

the outer distribution is given by

$$(5) \quad U_\theta(r) = Ar + \frac{B}{r}, \quad R_c \leq r \leq R_o$$

where U_θ is the tangential velocity, U_c – maximum core velocity, r = distance from vortex center, R_c – vortex core radius ($R_c = 0.5R_o$ is used), R_o – outer radius.

The constants A and B are chosen in such a way that the velocity matches at $r = R_c$ and decays to zero at $r = R_o$. The initial circulation Γ_0 is equal to $2\pi R_c U_c$. The vortex rotates counterclockwise.

An initial uniform density profile is taken inside the VDI ($\rho = \rho_l$ for $r \leq R_o$).

Results of computations for the case $D = 0.2$; $W = 1.0$ and Mach number of the shock wave $M = 1.5$ are presented in Fig 4. Density isolines are shown at the time moments: $50\mu s, 70\mu s, 90\mu s, 110\mu s, 157\mu s, 203\mu s, 421\mu s$. In order to visualize strengths of vortices, streamlines are drawn in Fig. 4e-h.

The induction of vorticity due to the interaction with a shock wave leads to stretching of the initial vortex in the y direction. The interaction of the initial and induced vorticity results in further distortion of the initial round vortex and its transformation to a "snake-like" structure with a long "tail" and a round "head" (Fig 4h).

Streamlines in the upper part of Figs. 4e-f are far from a typical pattern of flux streamlines near a single-point vortex. Later, initial and induced co-rotating counter-clockwise vorticity regions merge into one region of concentrated vorticity (see Figs. 4g-h). The second counter-clockwise vortex is weak and hardly seen on density isolines and streamlines. The streamlines become similar to those around a non-symmetrical vortex pair.

In order to examine the influence of governing parameters W and D on VDI patterns three cases are considered. The ratio Γ_1/Γ_0 is calculated for these cases and presented in Table 1.

TABLE 1
The VDI interaction with a propagating shock wave, Γ_0 initial circulation, Γ_1 induced circulation, Γ_0/Γ_{0A} , and Γ_1/Γ_{1A} are ratios of initial and induced circulation to those in the case A

case	$D = \rho_l/\rho_h$	$W = U_c/U_b$	Γ_0/Γ_{0A}	Γ_1/Γ_{1A}	Γ_0/Γ_1
A	0.2	1.0	1	1	0.888
B	0.2	0.5	0.5	1	0.444
C	0.5	1.0	1	0.4	1.776

In order to estimate the induced circulation due to interaction between a VDI and a shock wave, an approximate expression obtained for shock-bubble interaction [4] is used:

$$(6) \quad \Gamma_1 \approx \frac{4R_0}{V_s} \frac{\Delta p}{\rho_2} \frac{\rho_h - \rho_l}{\rho_l + \rho_h},$$

where ρ_2 and V_s are values of density and stream velocity behind the shock wave, Δp - pressure drop across the shock wave.

The above expression was derived under the assumption of a uniform density profile inside the bubble, i.e., the density gradient is only caused by the density difference between the bubble and ambiance.

However, the density profile inside the VDI is not uniform due to rotation, and we have to prove that the density gradient at the outward boundary of the VDI (at the inner side) is equal to zero.

After a short time (prior to the interaction) the profiles of pressure and density inside the VDI become close to those given by equilibrium equations for an isentropic steady vortex:

$$(7) \quad \frac{dp}{dr} = \frac{\rho U_\theta^2}{r},$$

$$(8) \quad \frac{p}{\rho^\gamma} = \frac{p_{amb}}{\rho_{amb}^\gamma},$$

with the boundary condition $p_1(r = R_o) = p_{amb}$, where the index 'amb' corresponds to the ambiance conditions either in front or behind the shock wave.

The radial density gradient is obtained from (8):

$$(9) \quad \frac{d\rho}{dr} \sim p^{1/\gamma-1} \frac{dp}{dr},$$

At the outward boundary of the VDI $U_\theta = 0$; therefore, the pressure gradient is equal to zero (see (7)). Thus, the density gradient is also equal to zero at the inner side of the VDI boundary.

Dimensionless coordinates $(x/R_o; y/R_o)$ of centers of vortices, relative density (ρ_c/ρ_{amb}) and pressure (p_c/p_{amb}) at these centers in $T = 421\mu s$ are presented in Table 2. Local pressure minimums are taken as centers of vortices.

TABLE 2

VDI shapes, $T = 421\mu s$, MV-main counter-clockwise vortex, SV1- clockwise vortex, SV2-second counter-clockwise vortex.

Case A

	x/R_o	y/R_o	ρ_c/ρ_{amb}	p_c/p_{amb}
MV	7.053	0.26	0.283	0.866
SV1	7.413	0.76	0.846	0.893
SV2	1.818	0.518	0.926	0.930

Case B

	x/R_o	y/R_o	ρ_c/ρ_{amb}	p_c/p_{amb}
MV	6.82	0.37	0.244	0.884
SV1	7.23	3.03	0.672	0.920
SV2	5.45	0.13	0.917	0.993

Case C

	x/R_o	y/R_o	ρ_c/ρ_{amb}	p_c/p_{amb}
MV	6.71	0.386	0.415	0.788
SV1	7.41	0.933	0.968	0.961
SV2	5.48	0.213	0.971	0.956

The overall strength of vortices characterized by $\Gamma_0 + \Gamma_1$ is maximum in case A and minimum in case C (see Table 1). Thus, the rotational velocity induced by vortices also reduces from case A to C. This rotational

velocity is added to the flux velocity U_b and influences the speed of the VDI. The x-coordinate of the main vortex is maximum in case A and minimum in case C.

Due to the process of the merge of two counter-clockwise vortices the y-coordinate of the main vortex is shifted up. This shift is most marked in case A.

Density values show that the core of the main vortex remains almost unmixed. The mixing is more complete for the clockwise vortex (SV1). The second counter-clockwise vortex (SV2) is mixed with ambiance and its density is close to the ambiance density.

The values of a local pressure minimum at the centers of vortices depend upon the rotational velocity and density inside the VDI (see (7)). The deeper pressure minimum in case C is explained by a 2.5-time higher initial density than that in cases A and B. In spite of a higher density level in case C, the pressure minimum in the center of SV1 vortex for case C is less than in cases A and B. This is due to a smaller induced vorticity (see (6)) and a corresponding lower level of rotational velocity.

Results for $W = 0.1, 0.25, 0.5, 1.0$ and the same initial density of the VDI $D = 0.2$, i.e., the computations are performed for a constant Γ_1 and various Γ_0 values, are presented. Density isolines for these cases are given in Fig 5a-d. The reduction of the initial rotational velocity leads to the formation of a more symmetrical shape of a rolled-up VDI. In the case of a small initial vortical velocity a nearly symmetrical vortex pair is obtained (Fig. 5a). For higher values of W the VDI structure becomes non-symmetrical. Pressure isolines are shown in Fig 5e-h and a pair of two round vortices appears for all values of W .

In all cases considered in this section the splitting of the VDI develops into a non-symmetrical vortex pair.

4. Interaction between VDI and expansion waves.

4.1. Estimation of induced circulation. The interaction between a Prandtl-Meyer expansion wave and light vortical inhomogeneity is considered (Fig. 3b).

In a Prandtl-Meyer expansion wave streamlines diverge around a sharp convex corner. The mechanism of inducing vorticity due to a non-parallel gradient of density at the outer boundary of the vortex and gradient of pressure is relevant for the case of VDI interaction with an expansion wave. However, in this case the VDI passes through the region of pressure gradient and not through a steep pressure drop as in the previous case of the interaction with a propagating shock wave. Therefore, the induced vorticity grows gradually and this is an essential difference between the two types of interaction.

A search is made for an appropriate estimation of the circulation induced in a VDI passing through an expansion wave. An expansion wave or a part of it can be characterized by upstream and downstream Mach numbers (M_u and M_d , respectively). The approximation (6) was derived under the assumptions of stepwise pressure drop across the shock wave and a constant velocity of a shock propagation. For the interaction between the VDI and an expansion wave an approximate expression for an induced circulation is written as a sum of the circulation induced due to interactions with n weak waves:

$$(10) \quad \Gamma_1 = \sum_{i=2}^n \left(\frac{4R}{V\rho} \right)_i (p_i - p_{i-1}) \left(\frac{\rho_h - \rho_l}{\rho_l + \rho_h} \right),$$

The variation of Mach number for all n waves is the same:¹

$$(11) \quad \Delta M = M_{i+1} - M_i = \frac{M_d - M_u}{n - 1}$$

¹Variation of local angle is not the same for these n waves

p_{i-1} and p_i denote upstream and downstream pressure of the $i - th$ wave, R_i radius of the VDI in the interaction with the $i - th$ wave, V_i is the gas velocity upstream of the $i - th$ wave and ρ_i is the gas density downstream of the $i - th$ wave.

The last term $(\rho_h - \rho_l)/(\rho_l + \rho_h)$ is constant due to the fact that the density ratio ρ_h/ρ_l remains the same for an isentropic expansion of gases with the same ratio of specific heats. The first term, $Q_i = R/(V\rho)_i$, is variable as the radius of the VDI, velocity and density of the mainstream vary through an expansion wave.

The normalized value of induced circulation can be written as:

$$(12) \quad \overline{\Gamma}_1 \approx \sum_{i=2}^n \frac{Q_i}{Q_d} \frac{(p_i - p_{i-1})}{P_d},$$

where $Q_d = (R/(V\rho))_d$ and P_d are values of the parameter Q and pressure downstream from the expansion wave.

The first term Q_i/Q_d in the above equation can be expressed as a function of Mach numbers M_i and M_d as follows. The radius of the VDI increases due to the expansion of streamlines. The distance between streamlines is inversely proportional to the mass flux, $V\rho$. Therefore, the above term is given by the following expression:

$$(13) \quad \frac{Q_i}{Q_d} = \left(\frac{V_d \rho_d}{V_i \rho_i} \right) \left(\frac{R_i}{R_d} \right) = \left(\frac{V_d \rho_d}{V_i \rho_i} \right)^2$$

Stream velocities are connected directly to Mach numbers:

$$(14) \quad \frac{V_d}{V_i} = \frac{M_d}{M_i} \frac{c_d}{c_i},$$

where c - speed of sound.

Substituting isentropic relations [13]

$$(15) \quad \frac{p_d}{p_i} = \left(\frac{1 + [(\gamma - 1)/2]M_i^2}{1 + [(\gamma - 1)/2]M_d^2} \right)^{\frac{\gamma}{\gamma-1}};$$

$$(16) \quad \frac{\rho_d}{\rho_i} = \left(\frac{1 + [(\gamma - 1)/2]M_i^2}{1 + [(\gamma - 1)/2]M_d^2} \right)^{\frac{1}{\gamma-1}};$$

$$(17) \quad \frac{c_d}{c_i} = \sqrt{\frac{T_d}{T_i}} = \left(\frac{1 + [(\gamma - 1)/2]M_i^2}{1 + [(\gamma - 1)/2]M_d^2} \right)^{\frac{1}{2}}$$

in (13), we have

$$(18) \quad \frac{Q_i}{Q_d} = \left(\frac{M_d}{M_i} \right)^2 \left(\frac{1 + [(\gamma - 1)/2]M_i^2}{1 + [(\gamma - 1)/2]M_d^2} \right)^{\frac{\gamma+1}{\gamma-1}}$$

The normalized induced circulation is computed by (12) vs. parameter $Z = M_i/M_d$ for $M_d = 1.25, 1.5, 1.75, 2.0$ and is shown in Fig 6a. Computations were performed in an "upstream" manner up to $M_u = 1$, i.e, for Z ranges from $1/M_d$ to 1. The number of n waves is taken equal to 500.

The computed induced circulation can be used for practical purposes. However, these numerical computations are somewhat cumbersome, especially if an inverse problem is to be solved, i.e., to find M_u or M_d for a required level of induced circulation. On the other hand, the expression (6) is also an approximation of circulation; thus, high accuracy of integration (10) is excessive.

If Q_i is replaced by its approximation Q_{app} , independent of the i -th wave, then Γ_1 is equal to Q_{app} multiplied by the difference of upstream and downstream pressures of the expansion wave. For this purpose the function $Q_i(Z)$ is studied in detail. Graphs of Q_i/Q_d vs the parameter Z are shown in Fig. 6b. One can see that it is a slightly convex function of Z and the first term in the sum (12) can be approximated by the average value of Q_u and Q_d :

$$(19) \quad \bar{\Gamma} \approx 0.5(Q_u + Q_d)(p_u - p_d)$$

The use of this approximation leads to an overestimation of induced circulation. A relative error between an approximate normalized induced circulation (19) and its value computed by (12) is shown in Fig. 7. The error increases with the M_d/M_u ratio and with M_d . For $Z > 0.75$, a relative error is no more than 10%.

A straightforward approximation of Γ in the same form as (6)

$$(20) \quad \bar{\Gamma} \approx \frac{R_u}{V_u \rho_d}(p_u - p_d)$$

leads to an underestimation of induced circulation. A relative error of the above approximation is more than that for the previous approximation (19)) (compare families of curves a and b in Fig. 7).

4.2. Numerical study . The VDI splitting after its interaction with an expansion wave is studied numerically. The expansion wave is considered provides the same amount of induced vorticity as for the previously studied shock wave-VDI interaction. The upstream Mach number is taken equal to the Mach number of the propagating shock in the previous case ($M_u = 1.5$). The value of induced circulation is computed by (6). We choose an angle θ of the convex corner such as to produce the desired Prandtl-Meyer expansion wave. For a given value of induced circulation the value of Q_d is obtained from the expression (19). The downstream Mach number M_d is calculated from equation (18). The angle of the convex corner is computed from

$$(21) \quad \theta = \nu(M_d) - \nu(M_u),$$

where $\nu(M)$ is the Prandtl-Meyer function [13]. For example, for $M_u = 1.5$ and the downstream Mach number $M_d = 2.01$ we obtain $\theta = 13^0$.

The distribution of rotational velocity is given by equations (4), (5). The maximum rotational velocity of the vortex is taken equal to $0.3M$. Thus, rotational velocity and initial circulation are the same as in the previous case of the VDI interaction with a propagating shock wave.

A numerical simulation of the interaction between the VDI and the Prandtl-Meyer expansion wave is performed as follows. First, the Euler equations are solved numerically to obtain a converged steady state for uniformed flow followed by a Prandtl-Meyer expansion wave. After this the VDI was introduced into the flowfield.

The computational domain and the VDI before the interaction are shown in Fig. 3b. Typical stages of the VDI interaction with the expansion wave and downstream are presented in Fig. 8. Similar patterns of the VDI shapes are observed after the interaction with shock and expansion waves. We recall that both

initial and induced circulations are equal in these two cases. However, the time scale is different in these cases. The effective time of VDI splitting after the interaction with the Prandtl-Meyer expansion wave can be computed by $T_{eff} = T - T_{exit}$, where T_{exit} is the time when the VDI just passes an expansion wave and get all induced vorticity from the interaction. Time should also be scaled by the mainstream velocity(2.01 M / 0.6 M = 3.35). The VDI shape after the interaction with a shock wave at time $T = 110\mu s$ (Fig. 4e) is compared with the VDI shape after interaction with the Prandtl-Meyer wave at the corresponding time moment $T_{eff} = 110 * 3.35 + 371 = 740\mu s$ (Fig. 8h). One can conclude that the vortical shapes are similar at corresponding times.

5. Conclusions. A numerical modeling of interaction between vortical density inhomogeneity (VDI) and shock and expansion waves was performed using a 2-D Euler code. The use of an adaptive unstructured numerical grid made it possible to obtain a high accuracy of computations and capture regions of induced vorticity for a moderate overall number of mesh points.

The structure of two counter-rotating non-symmetrical vortices was observed after the VDI interaction with a propagating shock wave.

Splitting of the VDI passing through the Prandtl-Meyer expansion wave was studied numerically. The main difference between the interaction of the VDI with a shock wave and an expansion fan is a gradual increase of the induced vorticity.

A method for the computation of an overall induced circulation as a result of the interaction with a number of weak waves was proposed. A simplified approximation for computing this value was derived and its accuracy was estimated analytically. This expression was used for the computation of the expansion angle of the Prandtl-Mayer wave for a given value of the induced circulation.

The obtained shapes of the VDI after the interaction with an expansion wave are close to those obtained for the shock-VDI interactions for the same values of initial and induced circulations. This gives us the opportunity to obtain VDI shapes in a complicated geometry by the simulation of the interaction with a straight propagating shock wave in a rectangular computational domain.

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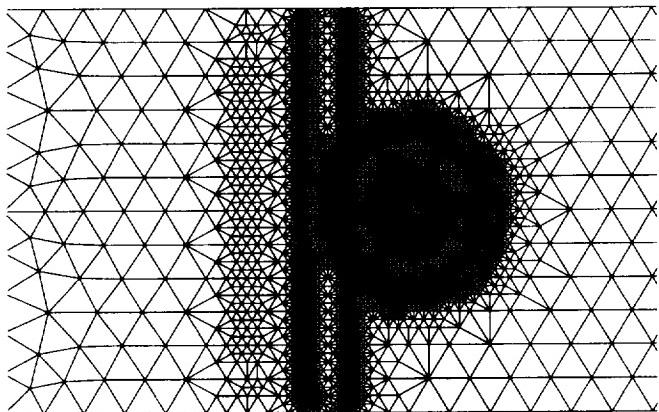
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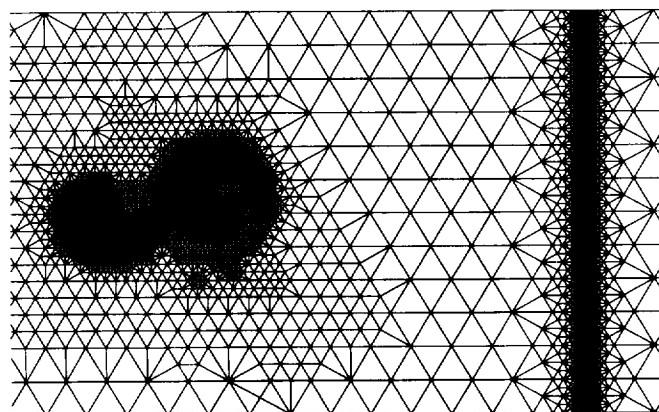
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a



b

FIG. 1. Fragment of unstructured grid with automatic mesh refinement: a, before the shock-VDI interaction; b, after the interaction

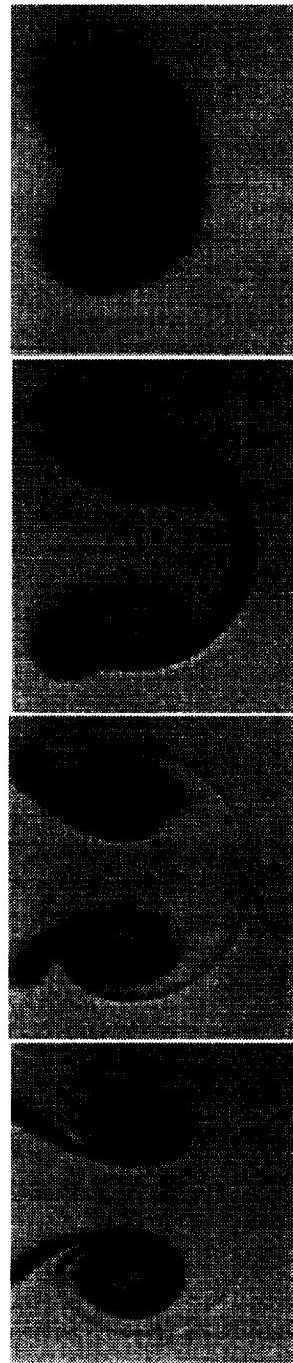
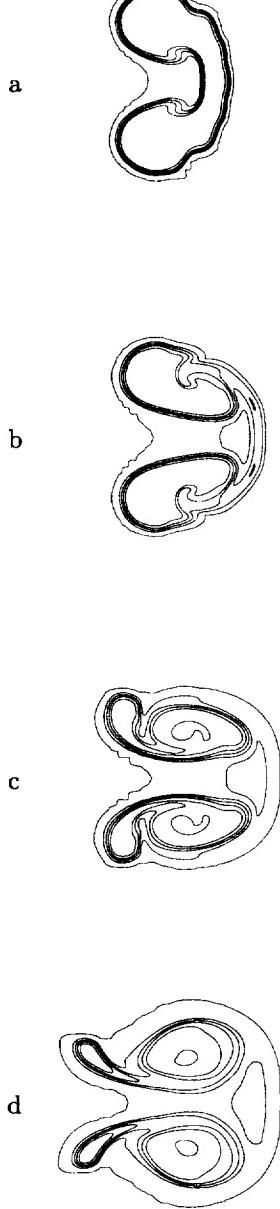
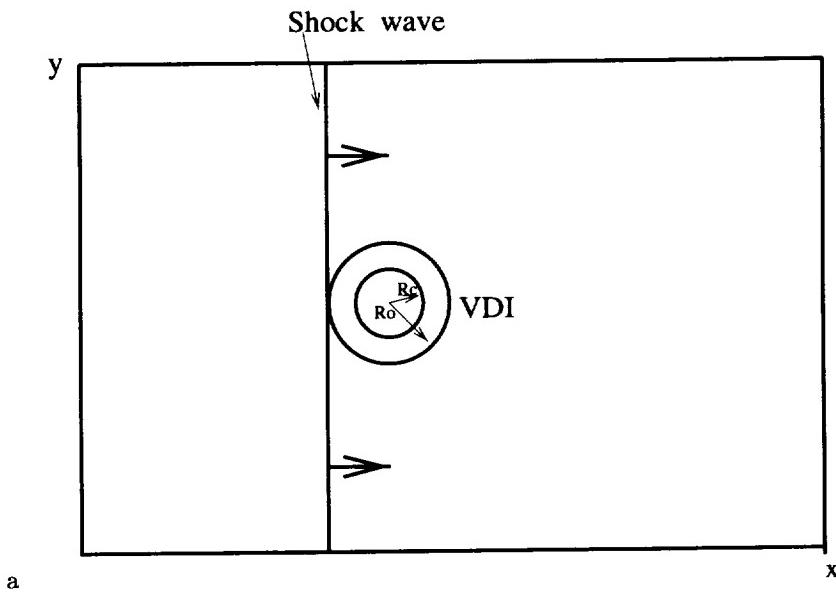
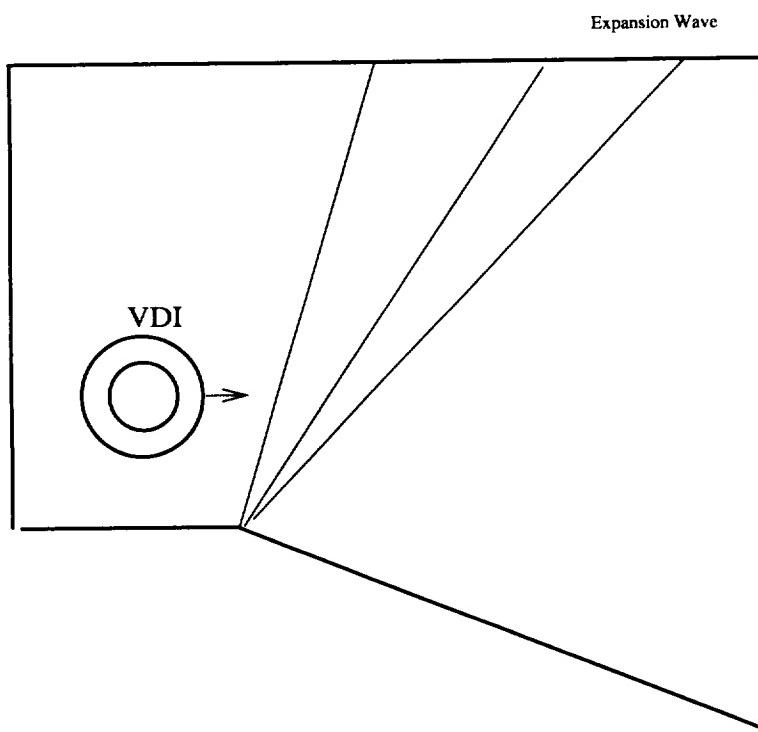


FIG. 2. Interaction of cylindrical bubble with a propagating shock wave ($M=1.09$; $D=0.14$; $W=0$). Density isolines are taken with 40 equally spaced intervals between the minimum density inside the bubble and ambient density. The left side-our computational results. The right side-experimental results [3]. Times: a, $123\mu s$; b, $373\mu s$; c, $573\mu s$; d, $773\mu s$



a



b

FIG. 3. Computational domain: a, interaction of the VDI with a propagating shock wave; b, interaction of the VDI with the Prandtl-Meyer expansion wave

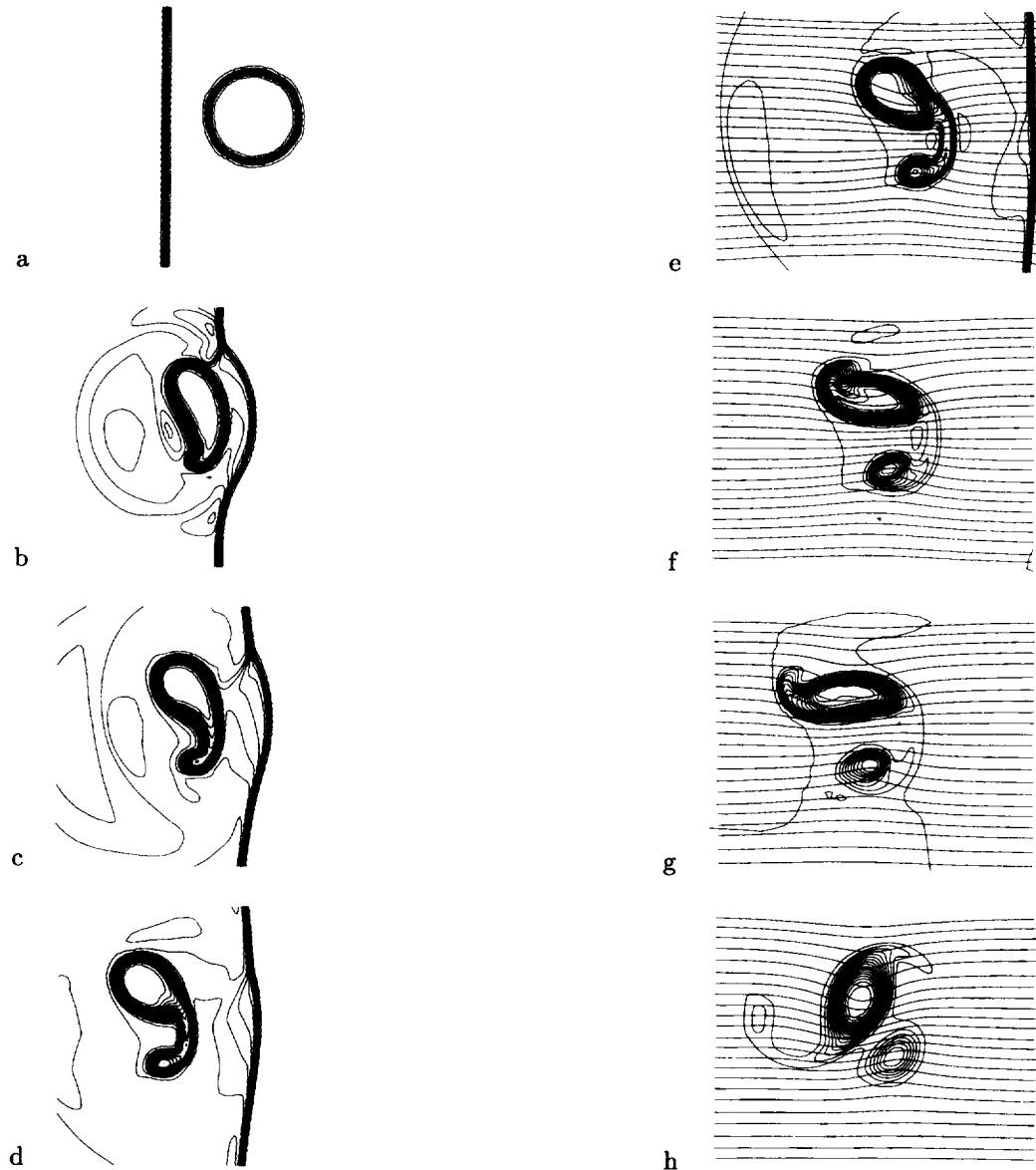


FIG. 4. Interaction of the VDI with a propagating shock wave ($D=0.2$; $W=0.5$) Density isolines are taken with 40 equally spaced intervals between density at the center of the counterclockwise vortex and ambient density. Times: a, 0; b, $50\mu s$; c, $70\mu s$; d, $90\mu s$; e, $110\mu s$; f, $157\mu s$; g, $203\mu s$; h, $421\mu s$

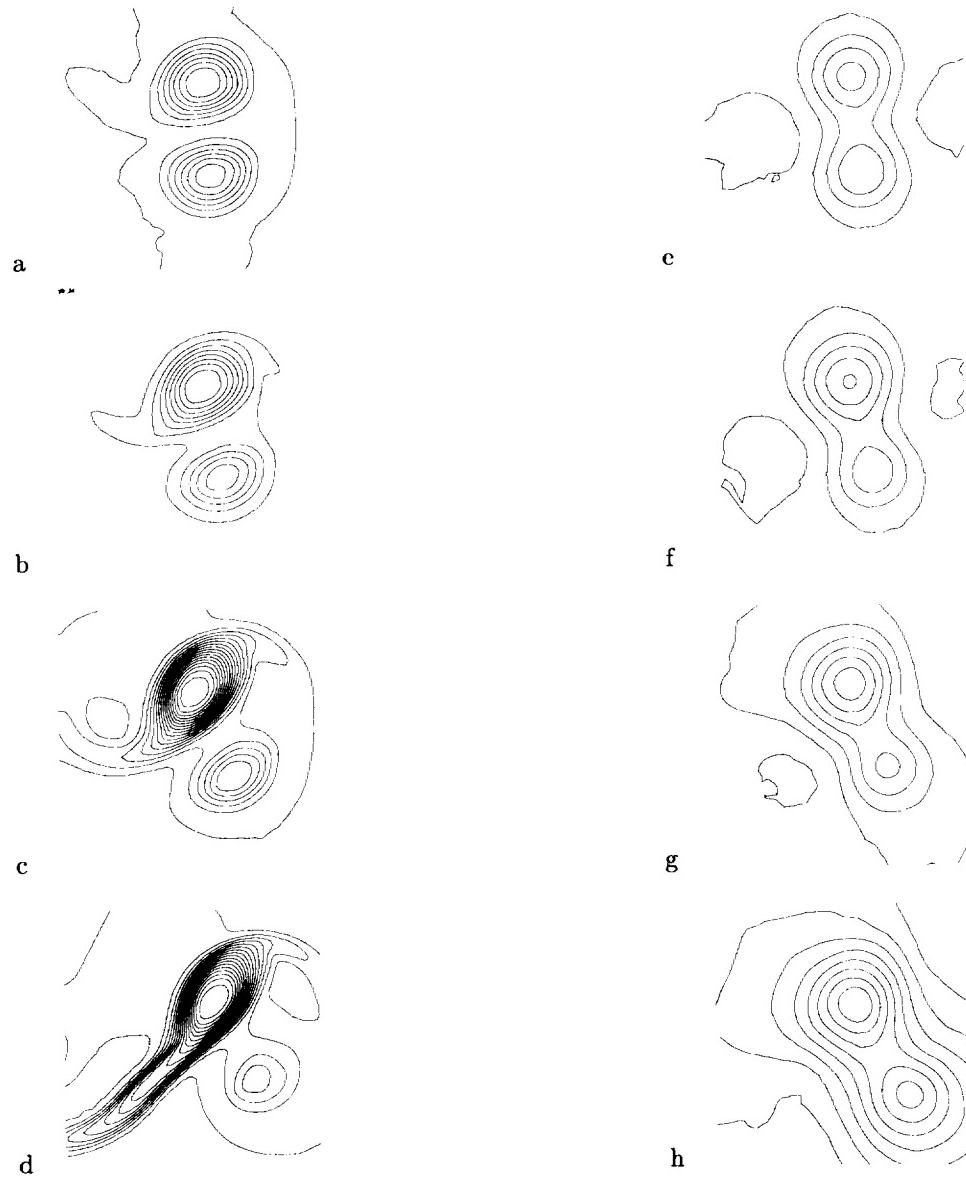


FIG. 5. Patterns of VDI shapes ($D = 0.2$) at $T = 421\mu s$: a-d, density isolines; e-h, pressure isolines, a,e, $W = 0.1$; b,f, $W = 0.25$; c,g, $W = 0.5$; d,h, $W = 1$

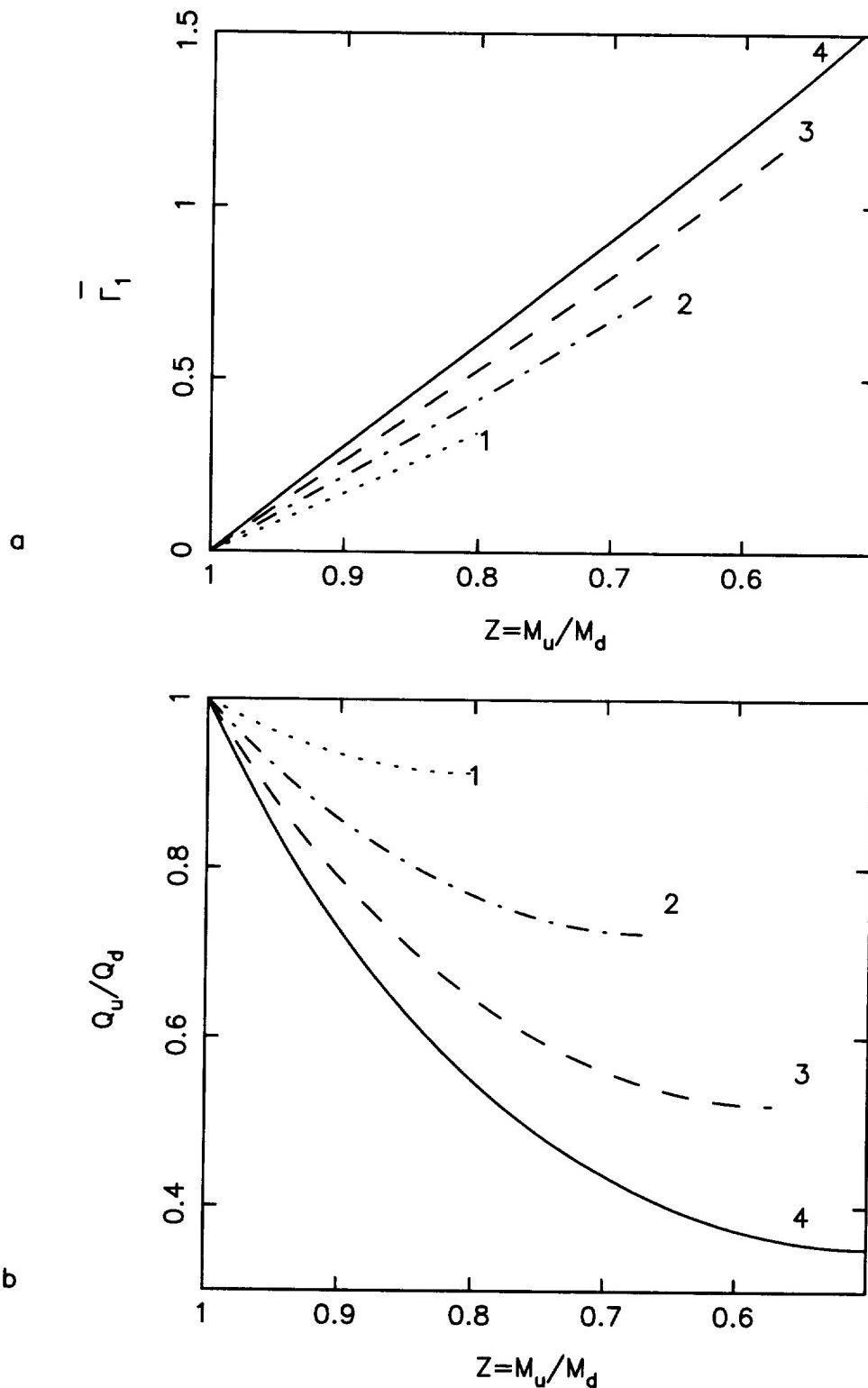


FIG. 6. Estimation of induced circulation due to the interaction between the VDI and the Prandtl-Meyer expansion wave:
 a, normalized induced circulation vs $Z = M_u/M_d$; b, normalized value of Q_u/Q_d vs Z . Downstream Mach numbers: 1- 1.25;
 2- 1.5; 3-1.75; 4-2.0

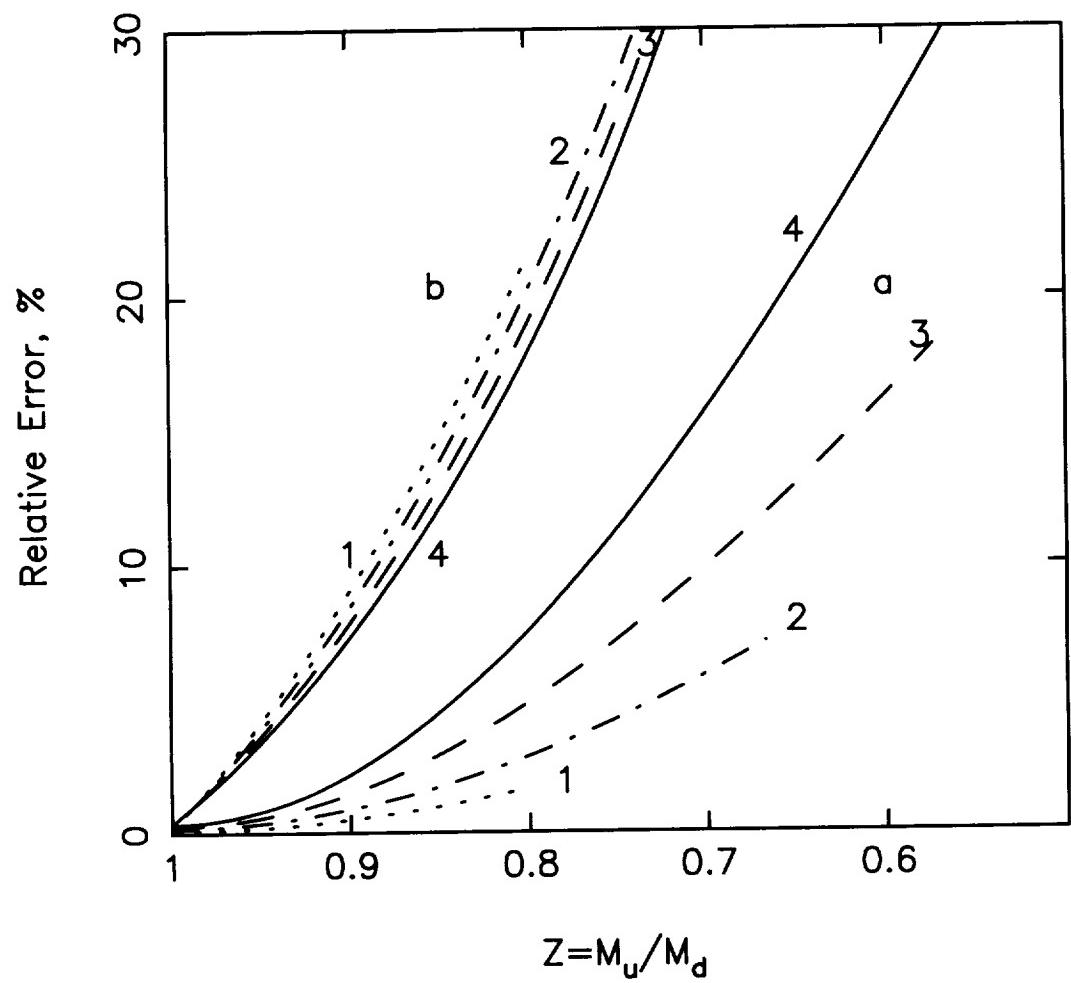


FIG. 7. Relative error of approximation for induced circulation: a, by (19); b, by (20). Downstream Mach numbers: 1-1.25; 2- 1.5; 3-1.75; 4-2.0

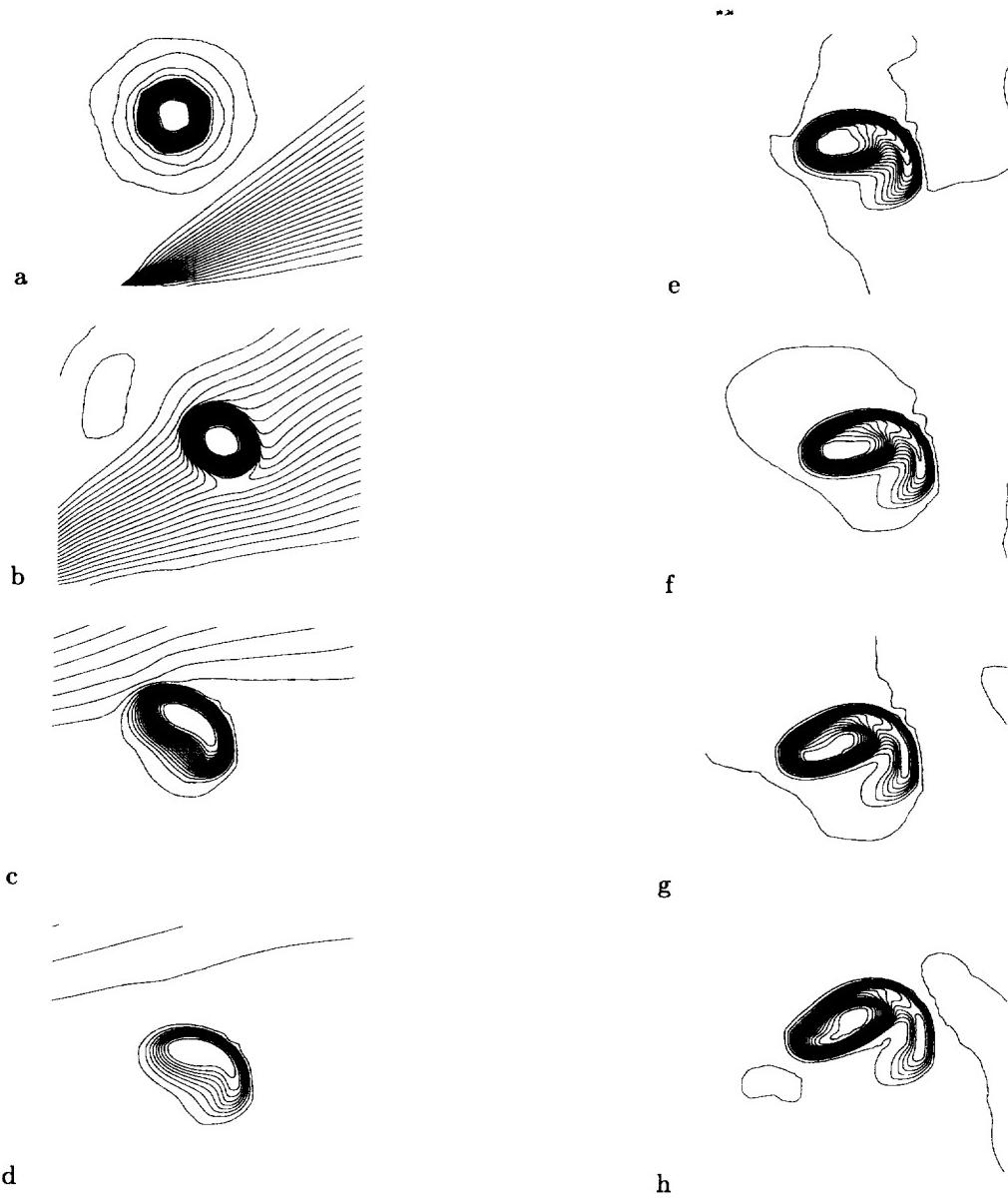


FIG. 8. Interaction of the VDI with the Prandtl-Meyer expansion wave ($D=0.2$). Density isolines are taken with 40 equally spaced intervals between minimum density inside the VDI and density upstream the expansion wave. Times: a, 0; b, $184\mu s$; c, $371\mu s$; d, $453\mu s$; e, $535\mu s$; f, $617\mu s$; g, $699\mu s$; h, $781\mu s$

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<p>We studied the interaction of a vortical density inhomogeneity (VDI) with shock and expansion waves. We call the VDI the region of concentrated vorticity (vortex) with a density different from that of ambiance. Non-parallel directions of the density gradient normal to the VDI surface and the pressure gradient across a shock wave results in an additional vorticity. The roll-up of the initial round VDI towards a non-symmetrical shape is studied numerically. Numerical modeling of this interaction is performed by a 2-D Euler code. The use of an adaptive unstructured numerical grid makes it possible to obtain high accuracy and capture regions of induced vorticity with a moderate overall number of mesh points. For the validation of the code, the computational results are compared with available experimental results and good agreement is obtained. The interaction of the VDI with a propagating shock wave is studied for a range of initial and induced circulations and obtained flow patterns are presented. The splitting of the VDI develops into the formation of a non-symmetrical vortex pair and not in a set of vortices. A method for the analytical computation of an overall induced circulation Γ_1 as a result of the interaction of a moving VDI with a number of waves is proposed. Simplified, approximated, expressions for Γ_1 are derived and their accuracy is discussed. The splitting of the VDI passing through the Prandtl-Meyer expansion wave is studied numerically. The obtained VDI patterns are compared to those for the interaction of the VDI with a propagating shock wave for the same values of initial and induced circulations. These patterns have similar shapes for corresponding time moments.</p>			
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